

## Frequency Response

Last time we

- Revisited formal definitions of linearity and time-invariance
- Found an eigenfunction for linear time-invariant systems
- Found the frequency response of a linear system to eigenfunction input
- Found the frequency response for cascade, feedback, difference equation, and differential equation systems

Today we will

- Extend the results to accommodate sinusoidal input, and then any input via Fourier series representation
- Write the Fourier series in terms of complex exponentials
- Provide a method to calculate Fourier series coefficients
- Determine properties of these coefficients

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## Response of an LTI System to Eigenfunction

Last time, we proved that for an input signal  $x$  given by

$$x(t) = e^{i\omega t} \quad \text{for all } t \in \text{Reals}$$

the corresponding output  $y$  of an LTI system can be expressed as

$$y(t) = H(\omega) e^{i\omega t} \quad \text{for all } t \in \text{Reals}$$

where  $H(\omega)$  is called the **frequency response** of the system.

The complex exponential  $e^{i\omega t}$  is called an **eigenfunction** of the system, because it creates an output with the same form, only differing by a scaling factor.

The same is true for a discrete input,

$$\begin{aligned} x(n) = e^{i\omega n} & \quad \text{for all } n \in \text{Integers} & \quad \text{leads to} \\ y(n) = H(\omega) e^{i\omega n} & \quad \text{for all } n \in \text{Integers} \end{aligned}$$

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## Cosines as Complex Exponentials

Recall that cosines can be expressed as complex exponentials:

$$\cos(\omega t) = \frac{e^{i\omega t} + e^{-i\omega t}}{2}$$

If we let  $x_1(t) = e^{i\omega t}$  and  $x_2(t) = e^{-i\omega t}$ , we express the cosine as

$$\cos(\omega t) = \frac{1}{2}(x_1(t) + x_2(t))$$

If we apply the cosine as input to an LTI system  $S$ , we find

$$S(\frac{1}{2}(x_1 + x_2)) = \frac{1}{2} (S(x_1) + S(x_2))$$

and since  $x_1$  and  $x_2$  are eigenfunctions, we can write

$$\begin{aligned} y(t) &= \frac{1}{2} (S(x_1) + S(x_2))(t) \\ &= \frac{1}{2} (H(\omega) e^{i\omega t} + H(-\omega) e^{-i\omega t}) \end{aligned}$$

So we can use frequency response to express the output for sinusoidal input.

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## Conjugate Symmetry

So for the input  $x(t) = \cos(\omega t)$ , we obtain the output

$$y(t) = \frac{1}{2} (H(\omega) e^{i\omega t} + H(-\omega) e^{-i\omega t}).$$

Realistic systems will produce purely real output (no imaginary component) for a purely real input like  $\cos(\omega t)$ .

This means that the imaginary parts of  $H(\omega) e^{i\omega t}$  and  $H(-\omega) e^{-i\omega t}$  must cancel out; they must be opposite in sign.

This is the same as saying that one is the conjugate of the other:

$$\begin{aligned} H(\omega) e^{i\omega t} &= (H(-\omega) e^{-i\omega t})^* \\ &= H(-\omega)^* e^{i\omega t} \end{aligned}$$

For systems that produce real output for real input, it is true that

$$H(\omega) = H(-\omega)^*$$

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## Implications: Scaled and Shifted Sinusoids

Systems that produce purely real output for a purely real input are called **conjugate symmetric**.

Let's look again at the output for our case  $x(t) = \cos(\omega t)$ ,

$$y(t) = \frac{1}{2} (H(\omega) e^{i\omega t} + H(-\omega) e^{-i\omega t})$$

Using the fact that  $z + z^* = 2 \operatorname{Re}\{z\}$ ,

$$\begin{aligned} y(t) &= \frac{1}{2} (2 \operatorname{Re}\{H(\omega) e^{i\omega t}\}) \\ &= \operatorname{Re}\{H(\omega) e^{i\omega t}\} \end{aligned}$$

If we express  $H(\omega)$  in polar form,  $H(\omega) = |H(\omega)| e^{i\angle H(\omega)}$ ,

$$\begin{aligned} y(t) &= \operatorname{Re}\{|H(\omega)| e^{i\angle H(\omega)} e^{i\omega t}\} \\ &= \operatorname{Re}\{|H(\omega)| e^{i(\angle H(\omega) + \omega t)}\} \\ &= |H(\omega)| \cos(\omega t + \angle H(\omega)) \end{aligned}$$

## Computing Sinusoidal Response

So, given the system response to an eigenfunction,  $H(\omega)$ , we can compute the **magnitude response**  $|H(\omega)|$  and the **phase response**  $\angle H(\omega)$ .

These form the scaling factor and phase shift in the output, respectively.

**The frequency of the output sinusoid will be the same as the frequency of the input sinusoid in any LTI system.**

**The LTI system scales and shifts sinusoids.**

These results hold true for both continuous and discrete signals and systems.



## Alternate Fourier Series Representation

Remembering that

$$\cos(\omega t) = \frac{e^{i\omega t} + e^{-i\omega t}}{2}$$

we may write

$$x(t) = A_0 + \sum_{k=1}^{\infty} \frac{A_k}{2} \left( e^{i(k\omega_0 t + \phi_k)} + e^{-i(k\omega_0 t + \phi_k)} \right) \quad \text{and also}$$

$$x(t) = A_0 + \sum_{k=1}^{\infty} \frac{A_k}{2} e^{i\phi_k} e^{ik\omega_0 t} + \frac{A_k}{2} e^{-i\phi_k} e^{-ik\omega_0 t} \quad \text{and letting}$$

$$X_k = \begin{cases} A_0 & \text{if } k = 0 \\ \frac{1}{2} A_k e^{i\phi_k} & \text{if } k > 0 \\ \frac{1}{2} A_{-k} e^{-i\phi_{-k}} & \text{if } k < 0 \end{cases}$$

$$x(t) = X_0 + \sum_{k=1}^{\infty} X_k e^{ik\omega_0 t} + X_{-k} e^{-ik\omega_0 t}$$

we obtain

## Alternate Fourier Series Representation: Discrete

For a discrete periodic signal, with the new notation

$$X_k = \begin{cases} A_0 & \text{if } k = 0 \\ \frac{1}{2} A_k e^{i\phi_k} & \text{if } k < \lfloor \frac{p}{2} \rfloor \\ A_k \cos(\phi_k) & \text{if } k = \lfloor \frac{p}{2} \rfloor \\ \frac{1}{2} A_{p-k} e^{-i\phi_{p-k}} & \text{if } k > \lfloor \frac{p}{2} \rfloor \end{cases}$$

$$x(n) = \sum_{k=0}^{p-1} X_k e^{ik\omega_0 n}$$

The proof is given in the text on page 303.

## Response to Fourier Series Input

Now let's apply a continuous input  $x(t)$  to an LTI system with frequency response  $H(\omega)$  and find the output  $y(t)$ :

$$x(t) = \sum_{k=-\infty}^{\infty} X_k e^{ik\omega_0 t}$$

Due to linearity, we can distribute over the sum and pull out the constants  $X_k$ .

The result is a scaled sum of the output generated by each individual complex exponential  $e^{ik\omega_0 t}$ .

Since each  $e^{i\omega t}$  has corresponding output  $H(\omega)e^{i\omega t}$ ,

$$y(t) = \sum_{k=-\infty}^{\infty} X_k H(k\omega_0) e^{ik\omega_0 t}$$

## Determining Fourier Series Coefficients

We now give formulae for the Fourier series coefficients for a periodic signal of period  $p$ :

For continuous signals  
 $x(t)$ ,  $t \in \text{Reals}$ :

$$X_m = \frac{1}{p} \int_0^p x(t) e^{-im\omega_0 t} dt$$

For discrete signals  $x(n)$   
 $x(n)$ ,  $n \in \text{Integers}$ :

$$X_m = \frac{1}{p} \sum_{n=0}^{p-1} x(n) e^{-im\omega_0 n}$$

The textbook provides a validation of these formulae on page 306, but their derivation will be intuitive once we have covered Fourier transforms.